

## Solution to Ex. 13.21

of *Turbulent Flows* by Stephen B. Pope, 2000

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April 13, 2017

For a general filter  $G(\mathbf{r}, \mathbf{x})$  satisfying the normalization condition Eq. (13.2), the filtered density function (Pope 1990) is defined by

$$\bar{f}(\mathbf{V}; \mathbf{x}, t) \equiv \int G(\mathbf{r}, \mathbf{x}) \delta[\mathbf{U}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{r} \quad (1)$$

Obtain the results

$$\bar{\mathbf{U}} = \int \mathbf{V} \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \quad (2)$$

$$\overline{U_i U_j} = \int V_i V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \quad (3)$$

$$\tau_{ij}^R = \int (V_i - \bar{U}_i)(V_j - \bar{U}_j) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \quad (4)$$

where integration is over all  $\mathbf{V}$ ; and  $\bar{\mathbf{U}}$ ,  $\overline{U_i U_j}$ , and  $\tau_{ij}^R$  are evaluated at  $\mathbf{x}$ ,  $t$ . Show that  $\bar{f}$  satisfies the normalization condition Eq. (12.1) and that, if the filter is everywhere non-negative, then  $\bar{f}$  is also non-negative, and hence has the properties of a joint PDF. Argue that, for such positive filters, the residual stress  $\tau_{ij}^R$  is positive semi-definite (Gao and O'Brien 1993). Use similar reasoning to show that  $L_{ij}^o$  and  $R_{ij}^o$  (Eqs. (13.100) and (13.102)) are also positive semi-definite for positive filters.

### Solution

$$\begin{aligned} \int \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} &\equiv \int \int G(\mathbf{r}, \mathbf{x}) \delta[\mathbf{U}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{r} d\mathbf{V} \\ &= \int G(\mathbf{r}, \mathbf{x}) \int \delta[\mathbf{U}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{V} d\mathbf{r} \\ &= \int G(\mathbf{r}, \mathbf{x}) d\mathbf{r} \\ &= 1 \end{aligned} \quad (5)$$

This means the filtered density function satisfies the normalization condition. If  $G(\mathbf{r}, \mathbf{x})$  is non-negative, from Eq. (1) it is clear that  $\bar{f}$  is also non-negative.

$$\begin{aligned}
\int \mathbf{V} \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} &= \int \mathbf{V} \int G(\mathbf{r}, \mathbf{x}) \delta[\mathbf{U}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{r} d\mathbf{V} \\
&= \int G(\mathbf{r}, \mathbf{x}) \int \mathbf{V} \delta[\mathbf{U}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{V} d\mathbf{r} \\
&= \int G(\mathbf{r}, \mathbf{x}) \mathbf{U}(\mathbf{x} - \mathbf{r}, t) d\mathbf{r} \\
&= \bar{\mathbf{U}}(\mathbf{x}, t)
\end{aligned} \tag{6}$$

$$\begin{aligned}
\int V_i V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} &= \int V_i V_j \int G(\mathbf{r}, \mathbf{x}) \delta[\mathbf{U}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{r} d\mathbf{V} \\
&= \int G(\mathbf{r}, \mathbf{x}) \int V_i V_j \delta[\mathbf{U}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{V} d\mathbf{r} \\
&= \int G(\mathbf{r}, \mathbf{x}) U_i(\mathbf{x} - \mathbf{r}, t) U_j(\mathbf{x} - \mathbf{r}, t) d\mathbf{r} \\
&= \overline{U_i U_j}(\mathbf{x}, t)
\end{aligned} \tag{7}$$

$$\begin{aligned}
&\int (V_i - \bar{U}_i)(V_j - \bar{U}_j) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \\
&= \int (V_i V_j - V_i \bar{U}_j - \bar{U}_i V_j + \bar{U}_i \bar{U}_j) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \\
&= \int V_i V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} - \bar{U}_j \int V_i \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} - \bar{U}_i \int V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} + \bar{U}_i \bar{U}_j \int \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \tag{8} \\
&= \overline{U_i U_j} - \bar{U}_j \bar{U}_i - \bar{U}_i \bar{U}_j + \bar{U}_i \bar{U}_j \\
&= \overline{U_i U_j} - \bar{U}_i \bar{U}_j \\
&= \tau_{ij}^R
\end{aligned}$$

From the above equations we can write

$$\begin{aligned}
\tau_{ij}^R &= \overline{U_i U_j} - \bar{U}_i \bar{U}_j \\
&= \int V_i V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} - \int V_i \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \int V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V}
\end{aligned} \tag{9}$$

For arbitrary vector  $\mathbf{Y}$ , and let  $A(\mathbf{V}) = Y_i V_i$ .

$$\begin{aligned}
\tau_{ij}^R Y_i Y_j &= Y_i Y_j (\overline{U_i U_j} - \bar{U}_i \bar{U}_j) \\
&= Y_i Y_j \int V_i V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} - Y_i Y_j \int V_i \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \int V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \\
&= \int Y_i Y_j V_i V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} - \int Y_i V_i \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \int Y_j V_j \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \\
&= \int A^2(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} - \int A(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \int A(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \\
&= \int \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \int A^2(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} - \int A(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \int A(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \\
&\geq^* \left( \int A(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \right)^2 - \int A(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \int A(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \\
&= 0
\end{aligned} \tag{10}$$

(10)

where the inequality marked by \* is the direct result of the Cauchy-Schwarz inequality<sup>1</sup> and the fact that  $\bar{f}$  is positive

$$\begin{aligned}
& \int \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \int A^2(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \\
&= \int \left[ \sqrt{\bar{f}(\mathbf{V}; \mathbf{x}, t)} \right]^2 d\mathbf{V} \int \left[ A(\mathbf{V}) \sqrt{\bar{f}(\mathbf{V}; \mathbf{x}, t)} \right]^2 d\mathbf{V} \\
&\geq \left[ \int \sqrt{\bar{f}(\mathbf{V}; \mathbf{x}, t)} A(\mathbf{V}) \sqrt{\bar{f}(\mathbf{V}; \mathbf{x}, t)} d\mathbf{V} \right]^2 \\
&= \int A(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} \int A(\mathbf{V}) \bar{f}(\mathbf{V}; \mathbf{x}, t) d\mathbf{V}
\end{aligned} \tag{11}$$

This derivation is inspired by my colleague DONG Bing <dongbing@sjtu.edu.cn>.

As for Leonard stresses  $L_{ij}^o$  we can define a new filtered density function

$$\bar{f}^L(\mathbf{V}; \mathbf{x}, t) \equiv \int G(\mathbf{r}, \mathbf{x}) \delta[\bar{\mathbf{U}}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{r} \tag{12}$$

Then

$$\begin{aligned}
\int \mathbf{V} \bar{f}^L(\mathbf{V}; \mathbf{x}, t) d\mathbf{V} &= \int \mathbf{V} \int G(\mathbf{r}, \mathbf{x}) \delta[\bar{\mathbf{U}}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{r} d\mathbf{V} \\
&= \int G(\mathbf{r}, \mathbf{x}) \int \mathbf{V} \delta[\bar{\mathbf{U}}(\mathbf{x} - \mathbf{r}, t) - \mathbf{V}] d\mathbf{V} d\mathbf{r} \\
&= \int G(\mathbf{r}, \mathbf{x}) \bar{\mathbf{U}}(\mathbf{x} - \mathbf{r}, t) d\mathbf{r} \\
&= \bar{\bar{\mathbf{U}}}(\mathbf{x}, t)
\end{aligned} \tag{13}$$

Here we have two pairs of analogies with  $\bar{f}^L$  to  $\bar{f}$  and  $\bar{\bar{\mathbf{U}}}$  to  $\bar{\mathbf{U}}$ . If we take  $\mathbf{U}$  to be a general random process  $\mathbf{B}$ , then Eq. (10) tells us that a tensor defined by  $\overline{B_i B_j} - \bar{B}_i \bar{B}_j = \overline{U_i U_j} - \bar{U}_i \bar{U}_j$  is positive semi-definite. Since  $L_{ij}^o$  is defined to be

$$L_{ij}^o = \overline{\bar{U}_i \bar{U}_j} - \bar{\bar{U}}_i \bar{\bar{U}}_j \tag{14}$$

And if we take  $\bar{\mathbf{U}}$  to be the general random process  $\mathbf{B}$ , then  $L_{ij}^o$  must be semi-definite due to the same reason expressed by Eq. (10).

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<sup>1</sup> [https://en.wikipedia.org/wiki/Cauchy%E2%80%93Schwarz\\_inequality](https://en.wikipedia.org/wiki/Cauchy%E2%80%93Schwarz_inequality)

Because the SGS Reynolds stresses are defined in the same way of  $L_{ij}^o$

$$R_{ij}^o = \overline{u'_i u'_j} - \overline{u'_i} \overline{u'_j} \quad (15)$$

And we take  $\mathbf{u}'$  as the general random process  $\mathbf{B}$ , then  $R_{ij}^o$  is also semi-definite.